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Electron energy band alignment at the (100)Si/MgO interface

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The electron energy band diagram at the (100)Si/MgO interface is characterized using internal photoemission of electrons and holes from Si into the oxide. For the as-deposited amorphous MgO the interface barrier corresponds to a band gap width of 6.1 eV, i.e., much lower than the conventionally assumed bulk crystal value (7.83 eV). The annealing-induced crystallization of MgO mostly affects the energy of the valence band while the conduction band bottom retains its energy position at 3.37 ± 0.05 eV above the top of the silicon valence band. © 2010 American Institute of Physics. [doi:10.1063/1.3294328]

Thanks to the ease of epitaxial growth on various crystal substrates1 and the wide bulk crystal band gap (7.83 eV)2 magnesium oxide, MgO, has attracted considerable interest as a material for a broad spectrum of electronic applications ranging from gate insulators in metal-oxide-semiconductor (MOS) transistors3–5 to a buffer layer enabling integration of other functional oxides.6,7 In particular, MgO is widely used as a barrier in magnetic tunnel junctions (MTJs)8–10 with the recently demonstrated possibility of spin injection into the silicon crystal.11,12 In all these applications, MgO is supposed to serve as an insulator providing high energy barriers for electrons and holes. Though the knowledge of the barriers is of particular importance for MTJs as they determine the type of charge carriers that provide the dominant contribution to the tunneling current,11 the band alignment between MgO and silicon, the most relevant for device fabrication, has not been addressed yet experimentally. One of the reasons for this lack of information concerns the strong impact of the x-ray induced insulator charging13 which in the case of MgO has been quoted as “commonly catastrophic.”14

To avoid the charging-related errors, here we have addressed the problem using the internal photoemission (IPE) spectroscopy.15–17 In IPE, the interface barrier height is obtained as the spectral threshold of photocurrent generated in a MOS capacitor by charge carriers photoexcited to energies above the barrier and subsequently injected into the insulating layer.15 Depending on the orientation of electric field at the interface, one may observe IPE of electrons or holes and determine the corresponding barrier heights.16,17 By combining the two barrier values, we have found that amorphous (a-) MgO on (100)Si has a much reduced band gap, Eg =6.1 eV, compared to the above mentioned bulk crystal value. Comparison between the as-deposited a-MgO layer and that crystallized by additional postdeposition anneal (PDA) indicates that the conduction band (CB) offset remains considerably lower (2.25 eV) than the valence band (VB) offset (2.7 eV in a-MgO and higher in the crystallized layer).

Studied samples were prepared on the H-terminated (100) surface of low-doped n- or p-type Si (nD, nA ~1015 cm−3) by e-beam deposition of MgO from 99.99% purity pellets with the Si temperature kept low (~150 °C) to obtain amorphous MgO layers.18 Some samples were subse-

FIG. 1. Logarithmic plot of the IPE quantum yield as a function of photon energy measured in as-deposited (a) (100)Si/MgO/Au samples and (b) after crystallization of the oxide insulators by PDA at 500 °C as measured under the indicated positive (open symbols) and negative (filled symbols) bias applied to the top Au electrode. Vertical lines mark the onsets of Eb and Ev of direct optical transitions A1–A1 and Σ2–Σ2, respectively, in the Si substrate crystal. Insets in panel (a) show schematics of the electron transitions.

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positive (open symbols) or negative (filled symbols) bias applied to the top Au electrode. In the photon energy range \( h\nu < 4.5 \) eV the curves measured under opposite orientation of electric field diverge suggesting that the photocurrent is dominated by IPE from the electrodes of the MOS capacitor. As replacement of the top electrode material from Au to Al has no effect on the curves measured under positive metal bias (not shown), electron states of the metal are unlikely to contribute to the IPE current thus indicating the electron IPE at the Si/MgO interface to be the major source. This conclusion is consistent with the observed modulation of the spectra at photon energies of \( E_1 = 3.4 \) eV and \( E_2 = 4.3 \) eV (cf. vertical lines in Fig. 1) due to the excitation of the direct optical transitions \( \Delta_{1}\Delta_{1} \) and \( \Sigma_{2}\Sigma_{2} \), respectively, in the Si crystal.\(^{16,17}\) These transitions cause a significant increase in reflectivity and optical absorption of Si but do not provide electrons of energy sufficient to overcome the interface barrier and, therefore, result in a decrease in the yield with increasing photon energy. The spectral threshold of electron IPE corresponds to the energy barrier between the top of the Si VB and the bottom of the MgO CB [cf. inset in the upper left corner in Fig. 1(a)] and can be found from the \( Y^{1/3} h\nu \) plots\(^{18,19}\) shown by the open symbols in Fig. 2 for the as-deposited (a) and crystallized (b) MgO. Though the electron IPE threshold is better exposed in the spectra of the crystalized layer [cf. Fig. 2(b)] than in those of a-MgO [panel (a)], the values obtained by linear extrapolation to the subthreshold level appear to be the same within experimental accuracy.

The IPE spectra measured under negative metal bias (filled symbols in Figs. 1 and 2) indicate [cf. panels (a)] that in the as-deposited sample the yield exhibits an enhancement at \( E_2 = 4.3 \) eV suggestive of hole IPE from Si.\(^{15}\) However, upon crystallization this feature disappears although a somewhat reduced photocurrent remains measurable in the range \( 3.5 < h\nu < 4.5 \) eV. This behavior can be explained by photocurrent arising in the same spectral range not only from hole IPE from the CB of Si into the MgO VB [cf. insert in the bottom right corner in Fig. 1(a)] but, also, from electron IPE from Au to the MgO CB since the Au/MgO interface is characterized by a barrier height of 4.0 eV.\(^{20}\) To separate electron and hole IPE signals it appears instructive to compare the yield spectra measured under the same bias of \( -7 \) V in samples with Au and Al electrodes prior and after MgO crystallization. The logarithmic plots of these spectra, shown in Fig. 3, reveal that crystallization of the MgO layer has no noticeable effect on the barrier height between the Fermi level of Al and the bottom of the MgO CB; the barrier is found to be \( 2.7 \pm 0.1 \) eV (the Fowler plots are not shown) in good agreement with the literature.\(^{20}\) By contrast, in the case of the Au electrode one may notice the disappearance upon annealing of the portion of current characterized by the well expressed \( E_2 \) singularity related feature. Therefore, we conclude that the crystallization of MgO suppresses the hole IPE from Si due to a shift of the MgO VB top. Since the hole IPE in the as-deposited samples provides the dominant photocurrent in the spectral range \( 4 < h\nu < 5 \) eV, for this case one still can evaluate the IPE spectral threshold from the \( Y^{1/3} h\nu \) plots\(^{18,19}\) as shown by the filled symbols in Fig. 2(a).

The field dependence of the electron IPE thresholds is shown in Fig. 4 in the form of Schottky plots for the amorphous (○) and crystallized (□) MgO. The measured electron IPE thresholds are seen to obey the classical image-force barrier model,\(^{15,21}\) yielding a zero-field barrier height of \( \Phi_0 = 3.37 \pm 0.05 \) eV and the high-frequency dielectric constant \( \varepsilon_{\text{MgO}} = 4.0 \pm 0.2. \) As one expects \( \varepsilon_1 \approx \varepsilon_2, \) the latter is in fair agreement with the measured refractive index \( n_{\text{MgO}} = 1.8 \) at \( h\nu = 5.0 \) eV. From the inferred energy barrier between the Si VB top and the MgO CB bottom the CB offset at the (100)Si/MgO interface can be calculated as \( \Delta E_{\text{C}} = \Phi \varepsilon_2 E_2(Si) = 0.22 \pm 0.005 \) eV.

The top of the image force barrier is located at a distance of \( x_m = (q/16\pi\varepsilon_0\varepsilon_2\varepsilon_{\text{MgO}} E_2)^{1/2} \) from the surface of emitter, where \( q \) is the elemental charge and \( \varepsilon_0 \) is the vacuum permittivity.\(^{15,17,21}\) As a result, the barrier height measured by IPE corresponds to the energy position of the MgO CB at the distance \( x_m \) from the silicon surface. In our experiments the strength of the electric field \( F \) in MgO increases from 0.1 to 2.5 MV/cm, which corresponds to a decrease in \( x_m \) from 3 to

*FIG. 2. Cube (left scale) and square root (right scale) plots of the IPE quantum yield as a function of photon energy measured in (100)Si/MgO/Au samples in the as-deposited state (a) and after MgO crystallization by PDA at 500 °C (b) under the indicated voltages applied to the Au electrode. Vertical lines mark the onsets of \( E_1 \) and \( E_2 \) direct optical transitions within the Si crystal and the spectral thresholds of electron (\( \Phi_e \)) and hole (\( \Phi_h \)) IPE.*

*FIG. 3. Logarithmic plot of the IPE quantum yield as a function of photon energy measured in (100)Si/MgO/Au samples with as-deposited (open symbols) and crystallized (filled symbols) MgO under \(-7\) V bias applied to the Au (circles) or Al (squares) electrode. The vertical line marks the onset of the \( E_2 \) direct optical transition in the Si substrate crystal.*

*FIG. 4. As-deposited and PDA 500 °C samples with Au or Al (circles) or Al (squares) electrode. The vertical line marks the onset of the \( E_2 \) direct optical transition in the Si substrate crystal.*
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