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# Supercritical Swelling of Liquid Crystal Films

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**Abstract**

The influence of liquid and supercritical carbon dioxide and liquid propane on the structural properties of both ionic and non-ionic surfactant-based liquid crystal films is discussed in this paper.  Swelling of the films, measured using *in-situ* small angle neutron scattering (SANS), was found to be dependent on the solubility of the propane/carbon dioxide in the micelles of the respective liquid crystals.  Additionally, under certain pressure conditions the structural properties of some of the films were

observed to change, ultimately leading to a loss of order in the micellar arrays of the liquid crystals.

**Introduction**

Investigation of polymeric liquid crystals (LCs) and their applications has become an important research area, with chemists exploiting the supramolecular ordering of LCs to create nanoscale materials. Of particular importance is LC templating of nanoporous materials, such as the mesoporous silicates MCM [1, 2], SBA [3] and MSU [4]. Ionic surfactants and non-ionic polymer LC templates are used in the synthesis of high quality mesostructured materials due to their supramolecular templating ability [5]. Many authors have reported the successful use of short chain ethylene oxide surfactants [6] and triblock copolymer surfactants containing poly(ethylene oxide) (PEO) and poly(propylene oxide) (PPO) segments [3, 7-9] in the synthesis of mesoporous materials.

LCs also have the ability to absorb significant amounts of solvent and consequently swell [10, 11]. However conventional swelling of LC templates with non-polar solvents requires careful consideration of many factors including the ionic strength of the solvent when swelling ionic surfactants, pH, temperature and the use a co-surfactant. To successfully swell the hexagonal phase of sodium dodecyl sulphate (SDS) Ramos *et al* [12] had to monitor the oil content of the LC solution and ionic strength of the swelling solvent and use a co-surfactant. In the study by Nagarajan *et al.* [13] swelling Pluronic aggregates with hydrocarbons resulted in cylinder to rod to lamellae micelle shape transitions.

Swollen liquid crystals (SLCs) have been utilized as templates for materials such as zirconia nanoneedles [14] and platinum nanoparticles [15] which were synthesised using sodium dodecyl sulfate SLCs, swollen with cyclohexane. In all cases the LC can easily be removed from the final material by destabilizing the SLC and washing with solvents. Work by Dos Santos and coworkers [16] has demonstrated the stability and dexterity of SLC templates over a wide range of pH values and compositions. Additionally, Yusuf *et al.* have utilized low molecular weight LCs to swell polymer thin films [11]. The researchers reported swelling of polymer thin films (approximately 150 μm thick) using low molecular weight LCs such as 4-*n*-pentyl-4-cyanobiphenyl and 4-methoxy-benzilidene-4-butyl-aniline [11]. They showed that upon swelling of films that their crystalline nature and structure was retained. Hanrahan *et al.* have demonstrated that supercritical carbon dioxide (sc-$CO_2$) can be utilised to swell LC templates during the production of mesoporous silica materials [17, 18].

Lubeck *et al.* used Brij 97 LCs as a template to produce CdS nanostructures [19] and Karanikolos *et al.* used a LC templated from P105 to direct the formation of ZnSe nanostructures [20]. By swelling the LCs, ZnSe nanomaterials with a variety of sizes could be produced. Alexandrisdis [21] *et al* demonstrated swelling of up to 20 % of Pluronic LCs templated form P105 using glycols.

In this paper we characterise the swelling of both ionic and non-ionic surfactants LC films with sc-$CO_2$ and high density propane. Structural changes of the micelles in the LCs were investigated using small angle neutron scattering (SANS) as a function of pressure and temperature. It was found that swelling (increasing the distance between hexagonal rods) with sc-$CO_2$ was easily achieved without any loss of ordering, *i.e.* no

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cylinder "wandering" was observed, as apposed to when conventional solvent swelling techniques are used. However, when propane was used to swell the same LCs problems such as structural rearrangements of the LCs at high propane pressures were encountered. The surfactants F127 [18], P123 [18], P85 [18], Brij 56 [22] and CTAB [23] at the concentrations used in this work have all been observed in the hexagonal phase in either LC experiments or as templates in mesoporous silica synthesis.

**Experimental**

*Materials and Reagents.* Hexadecyltrimethyl ammonium bromide (CTAB) (Fluka, UK) and F127 (PEO$_{100}$PPO$_{65}$PEO$_{100}$), P123 (PEO$_{20}$PPO$_{65}$PEO$_{20}$) and P85 (PEO$_{26}$PPO$_{40}$PEO$_{26}$), all polyethylene oxide-polypropylene oxide- polyethylene oxide triblock copolymers were used as received (BASF, New Jersey, U.S.A.). Deuterium (D$_2$O) and Brij 56 (CH$_3$(CH$_2$)$_{11}$EO ) were used as supplied by Sigma-Aldrich. All surfactants used in this work can be found in table 1.

*SANS Measurements of Liquid Crystal*. *In-situ* high pressure SANS experiments were conducted on the LOQ instrument at ISIS at the Rutherford Appleton Laboratory (RAL) in Oxfordshire, UK. The absolute scattering cross-section I (Q) (cm$^{-1}$) was measured as a function of the modulus of momentum transfer $Q(A^{-1}) = (4\pi/\lambda)\sin(\theta/2)$ where $\lambda$ is the incident neutron wavelength (2.2 to 10 Å) and $\theta$ is the scattering angle (< 7°). All measurements were conducted at 40 °C. High pressure SANS measurements were made using the UEA high pressure cell described in detail elsewhere [28]. Measurements performed on tertiary surfactant/D$_2$O/CO$_2$ mixtures were conducted at co-aqueous surfactant concentrations of 50 wt%.

A stock solution of LC, which consisted of 50 wt% surfactant (either F127, Brij 56 or CTAB) in $D_2O$, was prepared 2 hr before SANS measurements were undertaken. The surfactant was dissolved in the $D_2O$ while stirring at room temperature. The resultant viscous LC paste was applied evenly to the sapphire entrance window of the pressure cell. In order to prevent the dehydration of the LC film a pad of cotton wool soaked in $D_2O$ was inserted into the cell to saturate the $CO_2$ with water. The cell was allowed to equilibrate in both temperature and pressure for up to 30 min before SANS data was collected.

SANS data from the surfactant LC phase was consistent with the formation of a single pore size suggesting uniform dispersion of the copolymers within the cylindrical micelles which make up the hexagonal LC phases. The d-spacing (pore-centre-pore-centre distance, d) can be calculated using equation (1) [18] below:

$$d = \frac{4\pi}{Q_{MAX}\sqrt{3}} \tag{1}$$

where $Q_{Max}$ is the position of the diffraction peak maximum.

***Model for swelling of Liquid Crystals***. A simple model based on volume and surface area considerations was used to describe and understand the swelling of the LC phases. The model considers an initial volume of the LC phase (before swelling) filled with a number ($n_1$) of hexagonally arranged cylindrical micelles with fixed core radius $r_1$ and volume $V_{P1}$. This square slab is assumed to expand laterally but have fixed height. After swelling with $CO_2$ conservation of PPO:PEO interfacial area will dictate that the sample volume contains a reduced number of LC micelles ($n_2$) with

radius $r_2$ and PPO core (plus $CO_2$) volume $V_{P2}$. Equations (2) and (3) follow from this approach.

$$\frac{V_{P2}}{V_{P1}} = \frac{n_2 r_2^2}{n_1 r_1^2} \tag{2}$$

$$\frac{A_{S2}}{A_{S1}} = \frac{n_2 r_2}{n_1 r_1} \tag{3}$$

Where $A_{S1}$ and $A_{S2}$ are the mean area occupied by one surfactant molecule at the core-shell interface. Making the assumption that $A_{S1}=A_{S2}$, equations (2) and (3) yield equation (4) below:

$$\frac{V_{P2}}{V_{P1}} = \frac{r_2}{r_1} \tag{4}$$

The Pluronic LC d-spacing (D) changes through two effects, firstly an expansion in area of the slab as shown in equation (5),

$$d \propto a^{1/2} \propto (V_W + V_E + V_P)^{1/2} \tag{5}$$

Where $V_W$ = volume of water ($V_{W1}$ is before swelling and $V_{W2}$ is after swelling), $V_E$ = volume of PEO (it is assumed that the PEO micelle corona dimensions do not change), $V_P$ = volume of PPO ($V_{P1}$ before swelling and $V_{P2}$ after swelling)

Secondly, as a requirement that $A_{S1}=A_{S2}$, the number of micelle cylinders $n$ will decrease as $n \propto (1/r)$. This rearrangement will result in changes in an increase in d-spacing according to $d \propto (1/n)^{1/2}$. The overall result is that $d \propto r^{1/2}$ and since $r \propto V_P$ we obtain equation (6) below:

$$d \propto V_P^{1/2} \qquad \textbf{(6)}$$

The combined swelling given by (5) and (6) can be written as shown in equation (7),

$$\frac{d_2}{d_1} = \left(\frac{V_{W2} + V_E + V_{P2}}{V_{W1} + V_E + V_{P1}}\right)^{1/2} \left(\frac{V_{P2}}{V_{P1}}\right)^{1/2} \qquad \textbf{(7)}$$

Both the water (W) and PPO pseudo phases can be swollen by $CO_2$ and hence the above can be re-written in a form where the d-spacing increase with swelling (S) is given by,

$$\frac{d_S}{d} = \sqrt{S_{PO} \frac{V_{EO} + V_W S_W + V_{PO} S_{PO}}{V_{EO} + V_W + V_{PO}}} \qquad \textbf{(8)}$$

where $d_S$ represents the d-spacing of the swollen (S) phase and $d$ that of the non-swollen phase. This allows the swelling contributions of the PPO core and aqueous phase (W) to be expressed in terms of volume fractions $\phi_{CO_2-PPO}$ and $\phi_{CO_2-W}$ respectively as $S_{PO} = 1 + \phi_{CO_2-PO}$ and $S_W = 1 + \phi_{CO_2-W}$.

If expansion of interfacial area is allowed the above expression becomes (equation (9)):

$$\frac{d_S}{d} = \left(\frac{A_{S1}}{A_{S2}}\right)\sqrt{S_{PO}\,\frac{V_{EO} + V_W S_W + V_{PO} S_{PO}}{V_{EO} + V_W + V_{PO}}} \tag{9}$$

Here the maximum value of ($A_{S1}/A_{S2}$), corresponds to swelling with no increase in micelle number ($n_1 = n_2$) is $S_{PO}^{-0.5}$.

***Liquid Crystal Solubility measurements.*** $CO_2$ adsorption measurements were performed on a magnetic suspension balance, Rubotherm (Bochum, Germany). This method measures the change of weight of a sorbent due to adsorption of molecules from a surrounding compressed gas phase. The balance can be operated up to 500 bar using a magnetic coupling concept, which isolates the interior of the balance and the surrounding immediate environment. The weight gained by the substance is transmitted by magnetic suspension coupling from a closed and pressure-proof metal container to an external microbalance.

**Results and Discussion**

Figures 1 (a), (b) and (c) illustrate the effect of liquid $CO_2$ and sc-$CO_2$ pressure on the d-spacing of 50 wt% LC micelles of the Pluronic surfactant F127, the non-ionic surfactant Brij 56 and the ionic surfactant CTAB respectively as measured by SANS. The determination of structural arrangement of micelles was not the concern of this work and to identify the correct micelle packing would require further structural analysis of these systems. With increasing liquid $CO_2$ and sc-$CO_2$ pressure the d-spacing of the LC film micelles increases, as observed by the shift of $Q_{max}$ to lower Q values. Since the swelling of the aqueous phase is minimal ($S_W = 1 + \phi_{CO_2-W} \cong 1$) this increase in the d-spacing can be attributed solely to an increase in the micelle-to-micelle distance as the volume of the micelles swells with the uptake of liquid $CO_2$ and sc-$CO_2$, see figure 2. The corresponding increase in d-spacing obtained by application of equation 1 to the SANS data of figure 1 is shown as a function of sc-$CO_2$ pressure in figure 3.

The initial swelling of the F127 LC in both liquid $CO_2$ and sc-$CO_2$ shows a d-spacing increase from 175 Å to 203 Å upon increasing the pressure from 1 to 100 bar respectively, corresponding to an increase of 16 %. However, the increase in the d-spacing from 100 to 500 bar is very small, only 1 % of the overall swelling. The majority of the swelling takes place before 100 bar which is consistent with data reported for other Pluronic systems pressurised under sc-$CO_2$ [18]. This trend in the Pluronic LC swelling was also seen when the solubility of $CO_2$ in the 50 wt% LC of the Pluronic triblock copolymer P123 ($PEO_{20}PPO_{70}PEO_{20}$) was investigated by mass balance measurements, see figure 4. An increase in the solubility of $CO_2$ in the P123 LC is observed with increasing pressure. This increased $CO_2$ uptake is seen up to a

maximum at 100 bar, after which time the solubility is seen to level out. This result corresponds well with the SANS data obtained for the F127 LC and P123 LC data previously reported [18] and indicates that the solubility of the fluid in the surfactant causes the swelling.

The ability of $CO_2$ to swell the F127 micelle is due to the chemical composition of the micelle [24]. As with hydrocarbon solubilisation, $CO_2$ will predominantly dissolve in the polypropylene oxide (PPO) core of the micelles where the extent of solubilisation, and hence swelling of the LCs, increases with pressure.

The model used to obtain equation 8 links the d-spacing with the increase in micelle size as a function of $CO_2$ uptake. Using this approach, the radius of the Pluronic P123 LC micelle core was predicted for a number of pressures. These data are shown in figure 5 along with other calculated physical parameters of the P123 LC micelles under sc-$CO_2$ pressure. These data are presented as physical dimensions (Å) plotted against the PPO swelling factor ($S_{PPO}$) which is representative of the volume fraction of $CO_2$ taken up by the micelle core. Of course the micelle uptake of $CO_2$ depends on the pressure of $CO_2$. Also shown in figure 5 is the d-spacing and corresponding "wall thickness", *i.e.* the distance between corresponding micelle cores. Using pore size BJH calculations from calcined mesoporous silica synthesised using the P123 surfactant under sc-$CO_2$ from previous work [18], and taking into account calcination contraction ($\approx 30 \% $ [25]), the simple model predicts changes in dimensions of the $CO_2$-swollen LC phase of similar magnitude to those obtained from experimental measurements.

Since the $CO_2$ swells the PPO core of the micelles, the percentage PPO present in the liquid crystal micelle influences the increase in d-spacing. A greater extent of micelle expansion is seen in those Pluronic surfactants with a higher ratio of PPO to PEO. This is illustrated in figure 6 for the three Pluronics examined F127 (19 %) < P85 (25 %) < P123 (45 %) where the % of PPO to PEO blocks is given in brackets. However, application of the swelling model, with a conserved surfactant head group area, gives values for $\phi_{CO_2-PPO}$ as F127 (0.17) < P85 (0.25) < P123 (0.30). This would imply that $CO_2$ uptake by PPO declines with increasing EO number of the Pluronics. However, this is out of step with the general observation of a decreasing uptake with increasing molecular weight for liquids and polymers. An alternative explanation is that the head group area is not conserved but increases to a greater extent with increasing EO number. For this approach reasonable agreement is obtained for a constant $\phi_{CO_2-PPO} =$ 0.32 and the extent of head group expansion ($A_{S2}$ / $A_{S1}$) as F127 (1.15) < P85 (1.02) < P123 (1.0). Here F127 and P123 represent the extremes of maximum and minimum area expansion at the interface respectively.

Both the Brij 56 (figure 1b) and CTAB (figure 1c) LCs also show a significant shift of $Q_{max}$ with increasing sc-$CO_2$ pressure. Swelling of the Brij 56 and CTAB LCs resulted in an increase in the d-spacing of these films by 14.6 % and 5.6 % respectively.

The swelling of LCs with alkyl chain hydrophobes can be partly explained by considering the solubility of $CO_2$ in different n-alkanes. The uptake of sc-$CO_2$ has previously been reported to decrease as the length of the alkane chain increases [26]. However, as the Brij 56 ($C_{16}$) and CTAB ($C_{16}$) micellar cores have the same

composition, the extent of swelling should to be the same. However, when the structures of the CTAB and Brij 56 micelles are compared it is then reasonable to assume that the larger PEO units of the Brij 56 micelles affect the spacing between neighbouring micelles to such an extent that the distance between micelles is increased. This is mirrored in the swelling of Brij 56 and CTAB LCs by sc-$CO_2$. Interestingly, the swelling plateau observed for the F127 LC in sc-$CO_2$ is also seen in the Brij 56 and CTAB LCs; see figure 3.

The change in d-spacing at constant pressure (200 bar) with varying temperature for the F127 LC is shown in figure 7. At 200 bar the d-spacing is approximately 175 Å, which corresponds to the d-spacing shown in figure 3 of the F127 LC at atmospheric pressure. An increase in the d-spacing does not occur until the temperature approaches the critical temperature ($T_c$) of $CO_2$ (31.3 °C, the critical pressure of $CO_2$ is 74 bar). Once the $T_c$ of $CO_2$ is reached the viscosity and surface tension of the $CO_2$ atmosphere is reduced drastically and sc-$CO_2$ is able to penetrate into the highly viscous LC film, where the liquid $CO_2$ at low temperature was unable to do so.

As stated above the F127 LC is composed of a hydrophobic PPO micelle core and a hydrophilic PEO corona/interface [13, 24, 27]. It is the $CO_2$ – ether interactions that makes $CO_2$ so soluble in the PPO core [28] and it is the unique properties of sc-$CO_2$ [18] that improves its uptake of $CO_2$ into the micelles, as seen by inactive swelling below $T_C$. A degree of polarity is also important for $CO_2$ uptake, which is present in the PPO block but lacking in the PEO block [29]. The solubility of $CO_2$ in any polymer system can be explained by the plasticisation effect of $CO_2$ on the polymer. At increased pressure, gas molecules are forced between the polymer chains and expand the

spacing between the polymer chains which increases mobility. This results in more gas molecules being absorbed within the chains when the gas pressure is increased [30].

It is not thought that the "schizophrenic" nature, *i.e.* density fluctuations between $CO_2$ in the gas phase and $CO_2$ in the liquid phase and the tendency of $CO_2$ to prefer either gas or liquid at near critical conditions, is of concern in this work due to the fact that measurements were not conducted around the critical point of $CO_2$. Measurements were conducted above and below the critical point. A detailed swelling study around the critical points of $CO_2$ may illustrate anomalous swelling and/or $CO_2$ uptake due to density fluctuation domains.

Also shown in figure 7 is the F127 LC under the same temperature profile but in the absence of $CO_2$; no change in the micelle size was observed. Figure 7 also shows the d-spacing of a 50 wt% F127 LC at constant propane pressure with increasing temperature. However, the critical temperature of propane (96.7 °C) was not reached during this experiment and no similar change of d-spacing, as seen in the F127 liquid crystal in $CO_2$, was observed.

Figures 8 (a), (b) and (c) show how $Q_{max}$ for the 50 wt% Pluronic surfactants F127, the non-ionic surfactant Brij 56 and the ionic surfactant CTAB LCs change as a function of propane pressure, respectively. From the scattering pattern shown in figure 8 (a) it is clear that propane has no affect on the F127 LC micelles since there is no real change in the d-spacing (figure 9). King *et al.* [31] has previously reported that propane uptake decreases with increasing polymer size and since F127 is heavier than

those polymer surfactants studied, the lack of swelling can be attributed to the low solubility of propane in the relatively heavier tri-block copolymer F127 [31].

In contrast $Q_{max}$ was observed to decrease with increasing propane pressure for both the 50 wt% CTAB and Brij 56-templated LCs, indicating an increase in d-spacing. However, at high pressure both of these LC templates appear to be either destabilized or change morphology, as shown in figures 8 (b) and (c). In the case of Brij 56, $Q_{max}$ was observed to split into two separate peaks and lose its intensity above a propane pressure of 200 bar. However, when the pressure was reduced below 200 bar the peak was stable. This peak splitting may be the result of a change in morphology, possibly from hexagonally-ordered micellar cylinders to an unidentified arrangement of micelles, at high propane pressures.

Shown in figure 8 (c), prior to the destabilization of the CTAB LC by propane, is a very intense $Q_{Max}$ peak at approximately 0.084 Å$^{-1}$ (d-Spacing = 85 Å), up to a pressure of approximately 300 bar. Above this pressure $Q_{Max}$ drastically loses intensity and peak shape, but regains both when the pressure is returned to below 300 bar. Figure 9 shows the changes in d-spacing in F127, Brij 56 and CTAB LCs with increasing propane pressure before their destabilization.

The differences observed in the extent of swelling of the CTAB LCs when sc-$CO_2$ and propane gas were used as swelling agents, is indicative of the solubility of sc-$CO_2$ and propane gas in the CTAB LCs. Swelling the CTAB LCs at 300 bar in sc-$CO_2$ gives a 5.5 % increase in d-spacing, compared with an 11 % increase in the d-spacing at the same pressure for propane. The ability of propane to swell the CTAB LC to a

greater extent may be explained by propane's higher solubility in the long chain alkane core of the CTAB micelles.

This increased solubility of propane, compared to $CO_2$, in CTAB seems to be in agreement to that reported by Roy and co-workers [32], who investigated the solubility of propane and $CO_2$ in low concentrations of CTAB and sodium dodecyl sulphate (SDS) solutions.  At the CTAB concentrations used in this work the solubility of propane in CTAB would be greater than that of $CO_2$ in CTAB.

In the case of the Brij 56 LCs swelled by propane, it was shown by Beckman *et al.*, who studied the effect of liquid propane and ethane pressures on the structure of Brij 52 and Brij 30 microemulsions, that as the pressure of propane increases, the affinity of propane to the alkane tail of the Brij surfactants [33].  For this reason we see the initial swelling of Brij 56 by propane.

**Conclusions**

The data presented here demonstrates that surfactant LC films can be readily swollen in liquid carbon dioxide, supercritical carbon dioxide and liquid propane solvents. SANS data showed that the structural changes and swelling of the LCs was dependent on the solubility of the gases (propane or $CO_2$) in the respective micelles of the LCs. Modelling of the data showed that the micelle diameter increased with increasing $CO_2$ pressure, results in an increase in the observed d-spacing of the LCs.  Differences were also seen in the stability of the LCs in a propane gas environment compared to $CO_2$.  All the LCs were very stable in liquid $CO_2$ and sc-$CO_2$, however in a propane atmosphere some of them lost order or changed morphology, as observed by changes

in peak shape, peak splitting and a loss of intensity of $Q_{max}$. This work demonstrates the advantages of using sc-$CO_2$ over conventional swelling methods (hydrocarbon swelling) as the beneficial properties of $CO_2$ play an integral role in the swelling of the surfactant LCs.

**Acknowledgements**

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**Figures**

**Table 1**     The compositions of various surfactants used in this work

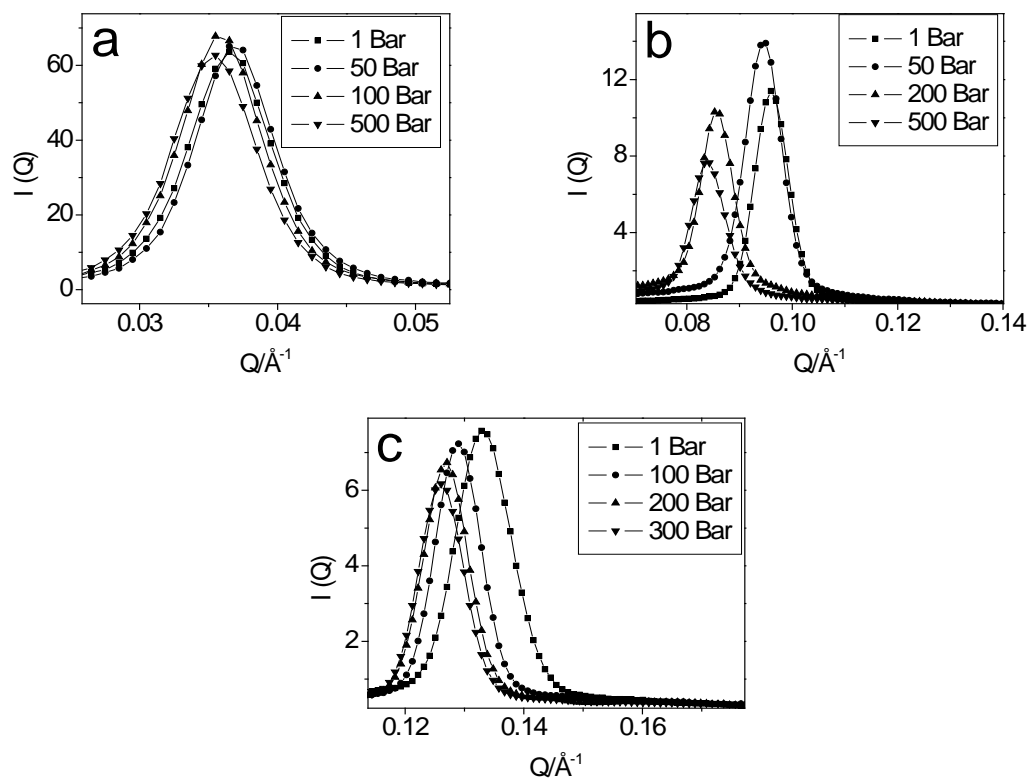| Surfactant | Composition |
|---|---|
| F127 | $PEO_{100}PPO_{70}PEO_{100}$ |
| P123 | $PEO_{20}PPO_{70}PEO_{20}$ |
| P85 | $PEO_{20}PPO_{40}PEO_{20}$ |
| CTAB | $(C_{16}H_{33})N(CH_3)_3$ Br |
| Brij 56 | $(C_{16}H_{33})$ $PEO_{10}$ |

**Figure 1**    SANS scattering patterns observed for different surfactants in a sc-$CO_2$ atmosphere as a function of pressure and at constant temperature of 40 °C.  (a), (b) and (c) show the effects of sc-$CO_2$ pressure on the F127, Brij 56 and CTAB LCs respectively.  Each figure has a different scale for illustration purposes.
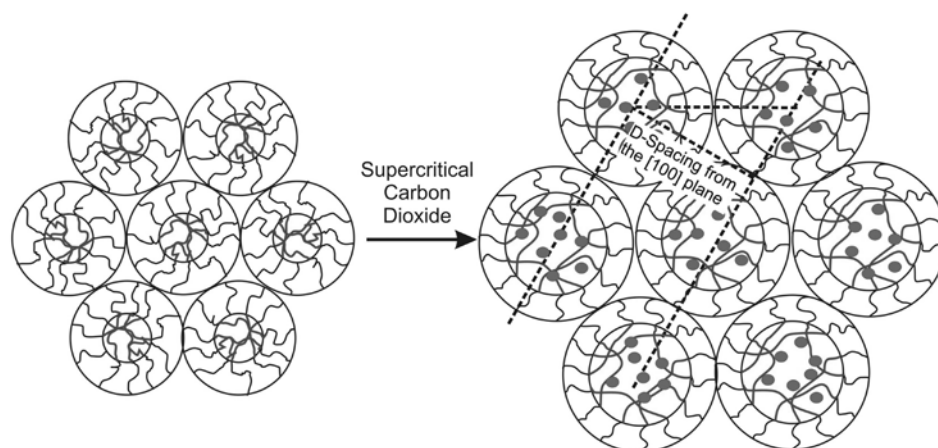
**Figure 2**  Schematic of the swelling process that takes place under an elevated sc-$CO_2$ pressure at a temperature of 40 °C.
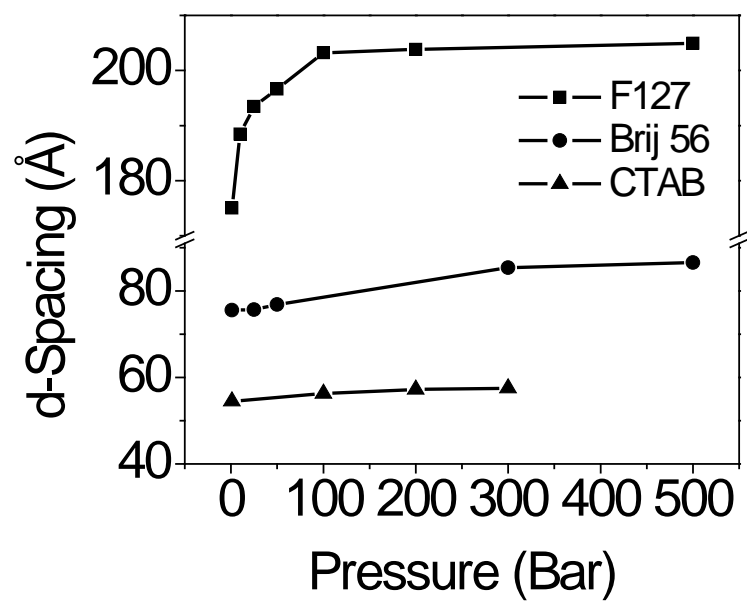
**Figure 3**    An overview of the structural changes observed in the LCs under a sc-$CO_2$ atmosphere as a function of pressure, at a constant temperature of 40 °C.  Figure shows the changes in d-spacings of F127 (■), Brij 56 (●) and CTAB (▲) LCs in sc-$CO_2$.
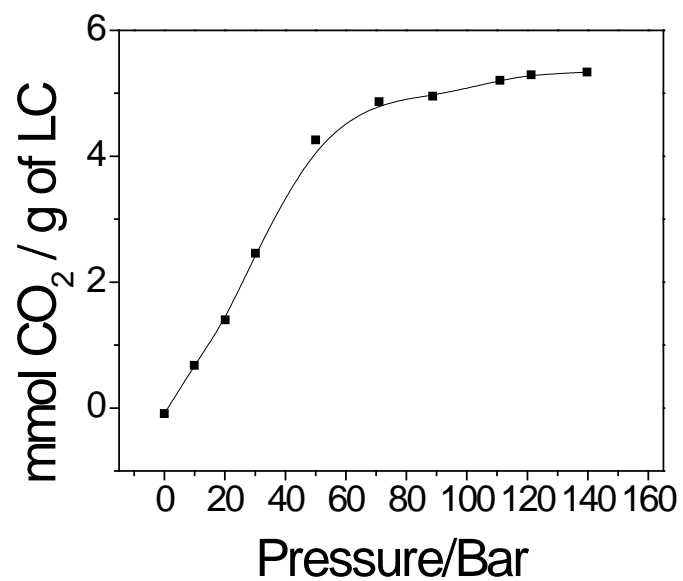
**Figure 4**     The adsorption of both liquid and sc-$CO_2$ in a P123 LC ($PEO_{20}PPO_{70}PEO_{20}$), determined by mass balance experiments conducted at a temperature of 40 °C.
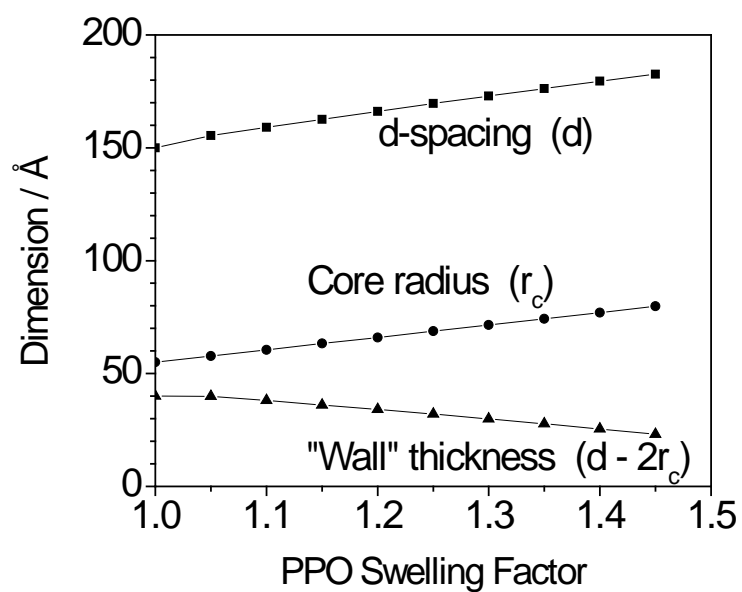
**Figure 5**     Physical parameters calculated from the SANS model (equation (9)) for the P123 LC micelles at a constant temperature of 40 °C. The PPO swelling factor scale depends on the size of the PPO micelle core.
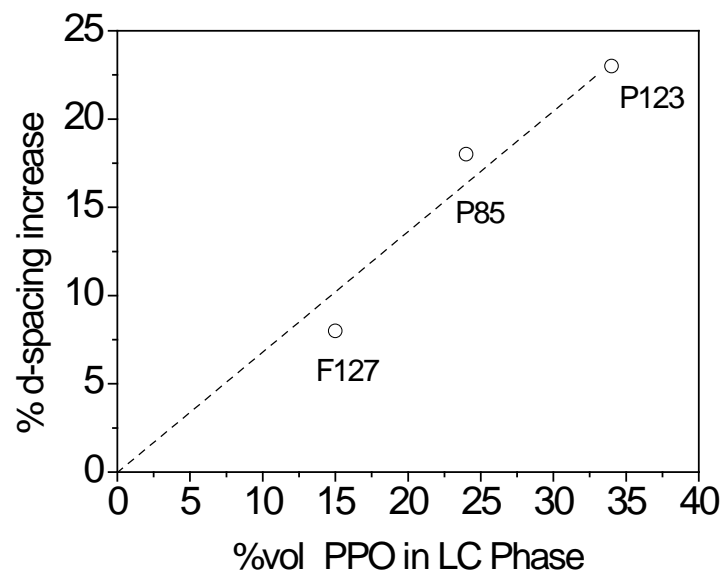
**Figure 6** Degree of swelling in sc-$CO_2$ at a pressure of 500 bar and a temperature of 40 °C. The composition of each Pluronic is P123: $PEO_{20}PPO_{70}PEO_{20}$, F127: $PEO_{100}PPO_{70}PEO_{100}$ and P85: $PEO_{26}PPO_{40}PEO_{26}$.
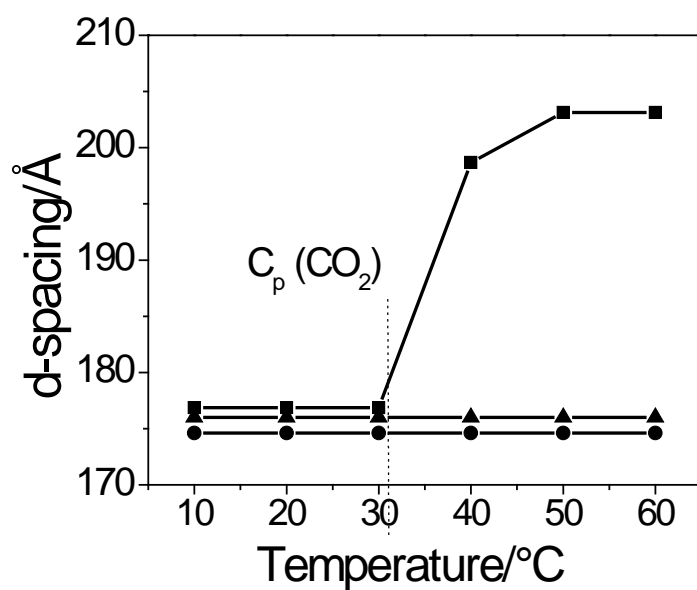
**Figure 7**   The influence of sc-$CO_2$ (■) and propane (●) on the d-spacing of the F127 LC as a function of temperature, at a constant pressure of 200 bar.  Also shown (▲) is the change in d-spacing of the F127 LC at atmospheric pressure.
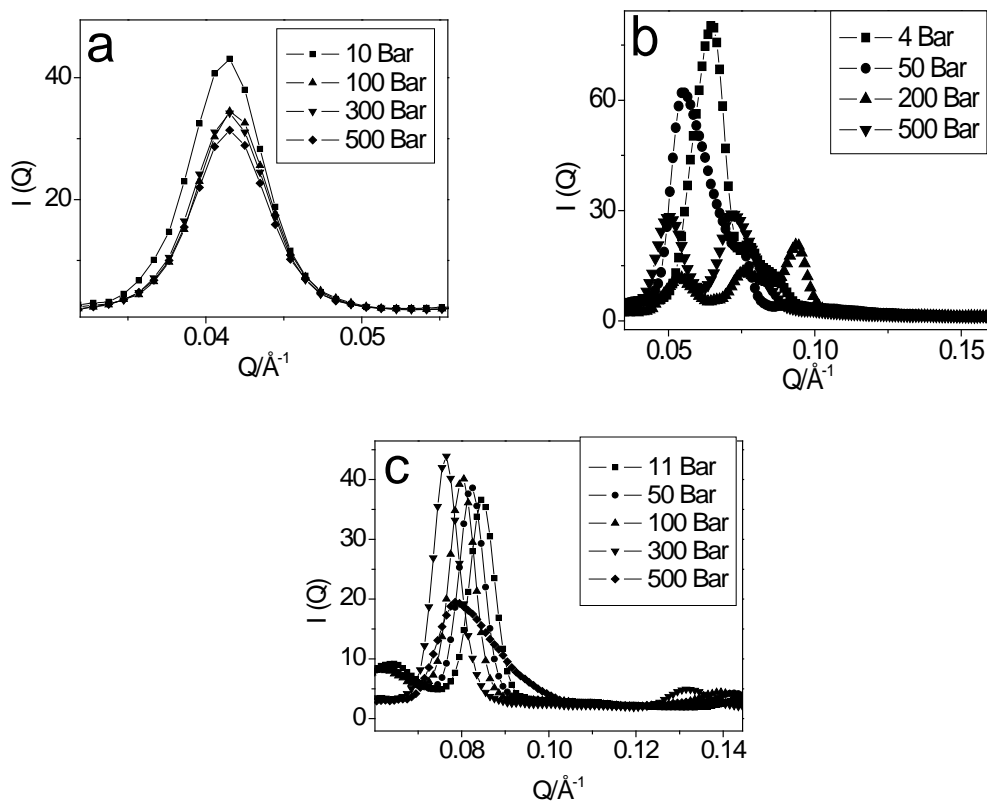
**Figure 8**    SANS scattering patterns observed for different surfactants in a propane atmosphere as a function of pressure, at a constant temperature of 40 °C. (a), (b) and (c) show the effects of propane pressure on the F127, Brij 56 and CTAB LCs respectively. Each figure has a different scale for illustration purposes.
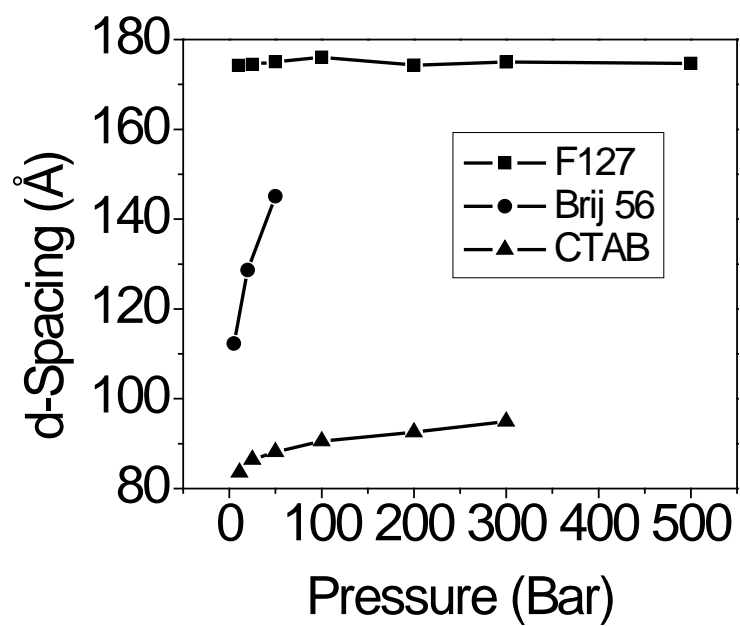
**Figure 9**    An overview of structural changes observed in LCs as a function of propane pressure, at a constant temperature of 40 °C. Figure shows the changes in d-spacings of F127 (■), Brij 56 (●) and CTAB (▲) LCs in propane.

# Swelling of Liquid Crystal Films

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